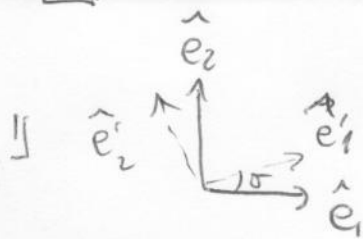


Feuille 2

II



$$x' = \cos\sigma x + \sin\sigma y$$

$$y' = -\sin\sigma x + \cos\sigma y$$

$$z' = z$$

rz

$$R_{ij} = \begin{pmatrix} \cos\sigma & \sin\sigma & 0 \\ -\sin\sigma & \cos\sigma & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

2) R_{ij} peut dépendre de tenseurs construits à partir de $n_i, \delta_{ij}, \epsilon_{ijk}$.

Pour faire un tenseur à 2 indices (de rang 2) :

$n_i n_j, \delta_{ij}, \epsilon_{ijk} n_k$.

$$R_{ij} = A \delta_{ij} + B n_i n_j + C \epsilon_{ijk} n_k$$

3) Pour une rotation autour de \hat{e}_3

$$R_{ij} = \begin{pmatrix} A & C & 0 \\ -C & A & 0 \\ 0 & 0 & A+B \end{pmatrix}$$

Par identification :

$$A = \cos\sigma$$

$$B = 1 - \cos\sigma$$

$$C = \sin\sigma$$

↳ Pour une rotation quelconque :

$$R_{ij} = \cos\sigma \delta_{ij} + (1 - \cos\sigma) n_i n_j + \sin\sigma \epsilon_{ijk} n_k$$

2

$$\begin{aligned}
 1) \quad [J_i, J_j]_{ke} &= (\bar{J}_i)_{km} (\bar{J}_j)_{me} - (\bar{J}_j)_{km} (\bar{J}_i)_{me} \\
 &= - \sum_{km} \epsilon_{kmi} \sum_{mej} + (i \leftrightarrow j) \\
 &= - (\delta_{ie} \delta_{jk} - \delta_{ij} \delta_{ke}) + (\delta_{je} \delta_{ik} - \delta_{ij} \delta_{ke}) \\
 &= \delta_{ik} \delta_{je} - \delta_{ie} \delta_{jk} \\
 &= \epsilon_{ijkm} \epsilon_{klem} \\
 &= i \epsilon_{ijm} (-i \epsilon_{klem})
 \end{aligned}$$

$$[\bar{J}_i, \bar{J}_j] = i \epsilon_{ijk} \bar{J}_k$$

$$\begin{aligned}
 2) \quad (\bar{J}_a n_a)_{ij}^2 &= (\bar{J}_a \bar{J}_b)_{ij} n_a n_b \\
 &= - \sum_{kha} \epsilon_{kja} \sum_{kib} \epsilon_{kib} n_a n_b \\
 &= (\delta_{ij} \delta_{ab} - \delta_{ib} \delta_{ja}) n_a n_b \\
 &= \delta_{ij} - n_i n_j
 \end{aligned}$$

$$\begin{aligned}
 (\bar{J}_a n_a)_{ij}^3 &= (\delta_{ik} - n_i n_k) (\bar{J}_a)_{kij} n_a \\
 &= -i (\delta_{ik} - n_i n_k) \epsilon_{kija} n_a \\
 &= -i \delta_{ija} n_a
 \end{aligned}$$

$$(\bar{J}_a n_a)^3 = \bar{J}_a n_a$$

$$(\mathbb{J}_a n_a)_{ij}^{2\alpha} = \delta_{ij} - n_i n_j \quad \forall \alpha > 0$$

$$(\mathbb{J}_a n_a)_{ij}^{2\alpha+1} = (\mathbb{J}_a n_a)_{ij} \quad \forall \alpha \geq 0$$

$$\begin{aligned} 3) \left[\exp(i\theta n_a \mathbb{J}_a) \right]_{ij} &= \sum_{\alpha} \frac{1}{\alpha!} (i\theta n_a \mathbb{J}_a)^{\alpha} \\ &= \delta_{ij} + \sum_{\alpha=1}^{\infty} \frac{(-1)^{\alpha}}{(2\alpha)!} \theta^{2\alpha} (\delta_{ij} - n_i n_j) + \sum_{\alpha=0}^{\infty} i \frac{(-1)^{\alpha}}{(2\alpha+1)!} \theta^{2\alpha+1} (n_a \mathbb{J}_a)_{ij} \\ &= \delta_{ij} + (\cos\theta - 1)(\delta_{ij} - n_i n_j) + i \sin\theta (n_a \mathbb{J}_a)_{ij} \\ &= \cos\theta \delta_{ij} + (1 - \cos\theta) n_i n_j + \sin\theta \varepsilon_{ija} n_a. \end{aligned}$$

Correspond à l'expression précédente...

$$\boxed{3} \quad \text{y} \quad \sigma_1^2 = \sigma_2^2 = \sigma_3^2 = \mathbb{1}.$$

$$\sigma_1 \sigma_2 = \begin{pmatrix} i & \\ & -i \end{pmatrix} = i \sigma_3$$

$$\sigma_2 \sigma_3 = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} = i \sigma_1$$

$$\sigma_3 \sigma_1 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = i \sigma_2.$$

$$(\sigma_i \sigma_j)^{\dagger} = \sigma_j^{\dagger} \sigma_i^{\dagger} = \sigma_j \sigma_i$$

$$\begin{aligned} \hookrightarrow \sigma_2 \sigma_1 &= -i \sigma_3 \\ \sigma_3 \sigma_2 &= -i \sigma_1 \\ \sigma_1 \sigma_3 &= -i \sigma_2. \end{aligned}$$

$$\hookrightarrow \sigma_i \sigma_j = \delta_{ij} \mathbb{1} + i \varepsilon_{ijk} \sigma_k.$$

$$2) \left[\frac{\sigma_i}{2}, \frac{\sigma_j}{2} \right] = \frac{\delta_{ij}}{4} \mathbb{1} + \frac{i}{4} \varepsilon_{ijk} \sigma_k - \frac{\delta_{ij}}{4} - \frac{i}{4} \varepsilon_{jik} \sigma_k = i \varepsilon_{ijk} \frac{\sigma_k}{2}.$$

$$3) (n_i \sigma_i)^2 = n_i n_j (\delta_{ij} \mathbb{1} + i \varepsilon_{ijk} \sigma_k) = \mathbb{1}$$

$$(n_i \sigma_i)^{2\alpha} = \mathbb{1} \quad \forall \alpha \geq 0$$

$$(n_i \sigma_i)^{2\alpha+1} = n_i \sigma_i \quad \forall \alpha \geq 0$$

$$\begin{aligned} 4) \exp\left(i \sigma n_i \frac{\sigma_i}{2}\right) &= \sum_{\alpha=0}^{\infty} (-1)^\alpha \left(\frac{\sigma}{2}\right)^{2\alpha} \mathbb{1} + \sum_{\alpha=0}^{\infty} i (-1)^\alpha \left(\frac{\sigma}{2}\right)^{2\alpha+1} n_i \sigma_i \\ &= \cos \frac{\sigma}{2} \mathbb{1} + i \sin \frac{\sigma}{2} n_i \sigma_i \end{aligned}$$

$$\begin{aligned} \boxed{4} \quad 1) \frac{1}{2} \text{Tr}(U U) &= \frac{1}{2} U_i V_j \text{Tr}(\sigma_i \sigma_j) \\ &= \frac{1}{2} U_i V_j \text{Tr}(\delta_{ij} \mathbb{1} + i \epsilon_{ijk} \sigma_k) \\ &= U_i V_i = \vec{U} \cdot \vec{V} \end{aligned}$$

$$\begin{aligned} -\frac{i}{2} \text{Tr}(U U W) &= -\frac{i}{2} U_i V_j W_k (\sigma_i \sigma_j \sigma_k) \\ &= -\frac{i}{2} U_i V_j W_k (\delta_{ij} \mathbb{1} + i \epsilon_{ijk} \sigma_k) \sigma_k \\ &= -\frac{i}{2} U_i V_j W_k (\delta_{ij} \sigma_k + i \epsilon_{ijk} \mathbb{1} - \epsilon_{ije} \epsilon_{ekm} \sigma_m) \\ &= \frac{1}{2} U_i V_j W_k \epsilon_{ijk} = (\vec{U} \times \vec{V}) \cdot \vec{W} \end{aligned}$$

$$\begin{aligned} 4) \text{Tr}(e_1' e_2' e_3') &= \text{Tr}(a^t e_1 a^t e_2 a^t e_3 a) \\ &= \text{Tr}(e_1 e_2 e_3) \end{aligned}$$

$$\hookrightarrow (\vec{e}_1 \times \vec{e}_2) \cdot \vec{e}_3 = (\vec{e}_1' \times \vec{e}_2') \cdot \vec{e}_3'$$

\hookrightarrow la transformation change une base directe en une base directe

\Rightarrow ce n'est pas une symétrie par rapport à un plan \Rightarrow se correspond à une rotation.

5) point de vue passif: le vecteur ne change pas:

$$\vec{U} = U_j \vec{e}_j = U'_j \vec{e}'_j$$

produit scalaire ~~\vec{e}_j~~ \vec{e}'_i

$$U_j \vec{e}_j \cdot \vec{e}'_i = U'_j \vec{e}'_j \cdot \vec{e}'_i = U'_i$$

$$\hookrightarrow U'_i = \underbrace{\vec{e}'_i \cdot \vec{e}_j}_{R_{ij}} U_j$$

$$6) R_{ij} = \vec{e}'_i \cdot \vec{e}_j = \frac{1}{2} \text{Tr} \left(\vec{e}'_i \vec{e}'_j \right) = \frac{1}{2} \text{Tr} \left(a^\dagger \vec{e}_i a \vec{e}_j \right)$$

$$= \frac{1}{2} \text{Tr} \left(a^\dagger \sigma_i a \sigma_j \right)$$

$$a^\dagger \sigma_i = \left(\omega \frac{\sigma}{2} \mathbb{1} - i \alpha_i \frac{\sigma}{2} n_a \sigma_a \right) \sigma_i$$

$$= \omega \frac{\sigma}{2} \sigma_i - i \sin \frac{\sigma}{2} \left(n_i \mathbb{1} + i n_a \epsilon_{aik} \sigma_k \right)$$

$a \sigma_j$: change $i \leftrightarrow j$
 $\sigma \leftrightarrow -\sigma$

$$a \sigma_j = \omega \frac{\sigma}{2} \sigma_j + i \sin \frac{\sigma}{2} \left(n_j \mathbb{1} + i n_a \epsilon_{rjk} \sigma_k \right)$$

$$a^\dagger \sigma_i a \sigma_j = \omega \frac{\sigma}{2}^2 \delta_{ij} + i \alpha_i \frac{\sigma}{2} \omega \frac{\sigma}{2} \left(i n_a \epsilon_{rji} - i n_a \epsilon_{oir} \right) +$$

$$n_a^2 \frac{\sigma}{2} \left(n_i n_j - n_a n_a \underbrace{\epsilon_{aik} \epsilon_{rjk}}_{(\delta_{or} \delta_{ij} - \delta_{or} \delta_{ri})} \right)$$

$$= \omega \frac{\sigma}{2}^2 \delta_{ij} + 2 \alpha_i \frac{\sigma}{2} \omega \frac{\sigma}{2} \epsilon_{ijk} n_k + \alpha_i^2 \frac{\sigma}{2} (n_i n_j - \delta_{ij})$$

$$= \left(\omega^2 \frac{\sigma}{2} - n_i^2 \frac{\sigma}{2} \right) \delta_{ij} + 2 n_i \frac{\sigma}{2} \omega \frac{\sigma}{2} \epsilon_{ijk} n_k + 2 n_i n_j n_l^2 \frac{\sigma}{2}$$

$$R_{ij} = \omega \sigma \delta_{ij} + (1 - \omega \sigma) n_i n_j + n_i \sigma \epsilon_{ijk} n_k$$

on retrouve toujours la même forme.

$$8] S_0^2 + S_1^2 + S_2^2 + S_3^2 = \omega^2 \frac{\sigma}{2} + n_i^2 \frac{\sigma}{2} \hat{n}^2 = 1$$

$$R_{ij}(\underline{S}) = (S_0^2 - S_d^2) \delta_{ij} + 2 S_0 S_k \epsilon_{ijk} + 2 S_i S_j$$

$R_{ij}(-\underline{S}) = R_{ij}(\underline{S})$: les points diamétralement opposés sont identifiés

$$a_{ij}(\underline{S}) = S_0 \delta_{ij} + i S_i \epsilon_j$$

$$a_{ij}(-\underline{S}) = - a_{ij}(\underline{S})$$

9]

S^2



RP^2



Si on combine 2 chemins



\approx



\approx



↕

chemin knoid.



$$\Rightarrow 1+1=0 \quad (\Rightarrow) \mathbb{Z}_2.$$

7

S_2 : tout chemin fermé se réfère au le chemin trivial.

10) à tout chemin fermé, on peut associer un nombre algébrique: le nombre de tours effectués des le sens trigonométrique. On peut ⁺ sommer les chemins \Rightarrow on somme la variable comptant le nombre de tours \Rightarrow \mathbb{Z} .

[5]

$$[p_i, p_j] = i\hbar \delta_{ij}$$

$$\begin{aligned} [e_i, n_j] &= \epsilon_{ike} [n_k p_e, n_j] = \epsilon_{ike} (n_k p_e n_j - n_j n_k p_e) \\ &= \epsilon_{ike} (n_k p_e n_j - n_k (p_e n_j + i\hbar \delta_{ej})) \\ &= -i\hbar \epsilon_{ikj} n_k = i\hbar \epsilon_{ijk} n_k \end{aligned}$$

$$\begin{aligned} [e_i, p_j] &= \epsilon_{ike} [n_k p_e, p_j] = \epsilon_{ike} (n_k p_e p_j - p_j n_k p_e) \\ &= \epsilon_{ike} (p_j n_k p_e + i\hbar \delta_{jk} p_e - p_j n_k p_e) \\ &= i\hbar \epsilon_{ijk} p_k \end{aligned}$$

$$\begin{aligned} [e_i, l_j] &= \epsilon_{ike} [e_i, n_k p_e] = \epsilon_{jke} (e_i n_k p_e - n_k p_e l_i) \\ &= \epsilon_{jke} (e_i n_k p_e - n_k l_i p_e + n_k l_i p_e - n_k p_e l_i) \\ &= \epsilon_{jke} ([e_i, n_k] p_e + n_k [e_i, p_e]) \\ &= i\hbar \epsilon_{jke} (\epsilon_{ikm} n_m p_e + \epsilon_{iem} n_k p_m) \\ &= i\hbar ((\delta_{em} \delta_{ij} - \delta_{ei} \delta_{mj}) n_m p_e + (\delta_{ik} \delta_{jm} - \delta_{ij} \delta_{mk}) n_k p_m) \\ &= i\hbar (n_i p_j - n_j p_i) \\ &= i\hbar \epsilon_{ijk} l_k \end{aligned}$$

$$\begin{aligned} [e_i, A_j^2] &= e_i A_j^2 - A_j e_i A_j + A_j e_i A_j - A_j^2 e_i \\ &= [e_i, A_j] A_j + A_j [e_i, A_j] \\ &= i\hbar (\epsilon_{ijk} A_k A_j + \epsilon_{ijk} A_j A_k) = 0 \end{aligned}$$

$$[e_i, f(A^2)] = 0.$$

Ces résultats nous montrent, comme on s'y attend que e_i est la génératrice des rotations (tous les vecteurs se transforment de la même façon, les scalaires ne se transforment pas).

$$\hookrightarrow [e_i, H] = 0.$$

\hookrightarrow Les valeurs propres de L^2 et L_z sont de bons nombres quantiques pour décrire un état propre de H .

$$2) \vec{R} = \vec{a} \times \vec{e} + \vec{v} \times \vec{e} - e^2 \frac{\vec{v}}{r} + e^2 \frac{\vec{p}}{r^2} \frac{dr}{dt}.$$

$= 0$ (potentiel central).

$$r = \sqrt{r_i^2}$$

$$\frac{dr}{dt} = \frac{\vec{p} \cdot \vec{v}}{r}$$

$$\vec{a} = \frac{F}{m} = -\frac{e^2}{m} \frac{\vec{p}}{r^3}$$

$$\begin{aligned} \vec{R} &= -\frac{e^2}{r^3} \vec{p} \times (\vec{p} \times \vec{v}) - e^2 \frac{\vec{v}}{r} + e^2 \frac{\vec{p}}{r^2} (\vec{p} \cdot \vec{v}) \\ &= -\frac{e^2}{r^3} \vec{p} (\vec{v} \cdot \vec{v}) + e^2 \frac{\vec{v}}{r} - e^2 \frac{\vec{v}}{r} + e^2 \frac{\vec{p}}{r^2} (\vec{p} \cdot \vec{v}) \\ &= \vec{0}. \end{aligned}$$

3) \hookrightarrow voir page suivante...

$$\begin{aligned} 4) R_i &= \frac{1}{2m} \epsilon_{ijk} (p_j p_k - l_j p_k) - e^2 \frac{p_i}{r} \\ &= \frac{1}{2m} \sum_{j,k} \epsilon_{ijk} (p_j p_k - p_k p_j - i\hbar \epsilon_{jke} p_e) - e^2 \frac{p_i}{r} \\ &= \frac{\epsilon_{ijk} p_j p_k}{m} - \frac{i\hbar}{m} p_i - e^2 \frac{p_i}{r}. \end{aligned}$$

$$3) R_i^* = \frac{1}{2m} \sum_{ijk} \epsilon_{ijk} (p_j p_k - p_k p_j) - e^i \frac{n_i}{n}$$

$$R_i^{\dagger} = \frac{1}{2m} \sum_{ijk} \epsilon_{ijk} (e^{\dagger}_k p_j^{\dagger} - p_k^{\dagger} e_j^{\dagger}) - e^i \frac{n_i^{\dagger}}{n^{\dagger}}$$

$$= \frac{1}{2m} \sum_{ijk} \epsilon_{ijk} (p_k p_j - p_k p_j) - e^i \frac{n_i}{n}$$

$$= + \frac{1}{2m} \sum_{ijk} \epsilon_{ijk} (p_k p_j - p_k p_j) - e^i \frac{n_i}{n}$$

$$= R_i$$

$$L_i R_i = \frac{\epsilon_{ijk}}{m} p_j p_k - \frac{i\hbar}{m} p_i - \frac{e^2}{a} p_i r_i$$

$$= \frac{\epsilon_{ijk}}{m} (p_j p_k + i\hbar \epsilon_{ise} p_e p_k) - \frac{i\hbar}{m} p_i p_i - \frac{e^2}{a} p_i r_i$$

$$\epsilon_{ijk} p_i p_j = 0$$

$$L_i p_i = \epsilon_{ijk} r_j p_k p_i = 0$$

$$p_i L_i = p_i p_i + \underbrace{[p_i, L_i]}_{=0}$$

$$L_i r_i = r_i L_i = r_i \epsilon_{ijk} r_j p_k = 0$$

$$\hookrightarrow \vec{L} \cdot \vec{R} = 0$$

$$\text{5)} [R_i, H] = \left[\frac{\epsilon_{ijk}}{m} p_j p_k - \frac{i\hbar}{m} p_i - e^2 \frac{r_i}{a}, \frac{p^2}{2m} - \frac{e^2}{a} \right]$$

$$[A, D] = \frac{\epsilon_{ijk}}{2m^2} \left(p_j p_k p^2 - p_j p^2 p_k + p_j p^2 p_k - p^2 p_j p_k \right)$$

$$[A, D] \stackrel{=0}{=} p_j \underbrace{[p_k, p^2]}_{=0} + \underbrace{[p_j, p^2]}_{=0} p_k$$

$$[A, E] = -\frac{e^2}{m} \epsilon_{ijk} \left(p_j p_k \frac{1}{a} - p_j \frac{1}{a} p_k + p_j \frac{1}{a} p_k - \frac{1}{a} p_j p_k \right)$$

$$= p_j \underbrace{[p_k, \frac{1}{a}]}_{=0} + \underbrace{[p_j, \frac{1}{a}]}_{=0} p_k$$

$$[p_i, r_j^2] = p_i r_j^2 - r_j^2 p_i + r_j p_i r_j - r_j^2 p_i = [p_i, r_j^2] r_j + r_j [p_i, r_j^2]$$

$$= -2i\hbar r_j r_i$$

$$[P_i, f(r^2)] = -2i\hbar r_i f'(r^2).$$

en particulier:

$$\begin{aligned} [P_i, \frac{1}{r}] &= -2i\hbar r_i \left(-\frac{1}{r}\right) \frac{1}{r^3} \\ &= i\hbar \frac{r_i}{r^3} \end{aligned}$$

$$[P_i, \frac{1}{r^3}] = 3i\hbar \frac{r_i}{r^5}$$

⋮

$$\begin{aligned} [A, E] &= -\frac{e^2}{m} i\hbar \frac{1}{r^3} \sum_{jk} \pi_j p_k \\ &= -\frac{i\hbar e^2}{m} \frac{1}{r^3} \underbrace{\left(\sum_{jk} \epsilon_{ijk} \epsilon_{kjm} r_j r_e p_m \right)}_{-r^2 p_i + r_i (\vec{r} \cdot \vec{p})} \end{aligned}$$

$$[A, E] = +\frac{i\hbar e^2}{m} \left(\frac{1}{r} p_i - \frac{r_i}{r^3} (\vec{r} \cdot \vec{p}) \right)$$

$$[B, D] = 0$$

$$\begin{aligned} [B, E] &= \frac{i\hbar e^2}{m} \left(p_i \frac{1}{r} - \frac{1}{r} p_i \right) \\ &= \frac{i\hbar e^2}{m} \left(\frac{1}{r} p_i + i\hbar \frac{r_i}{r^3} - \frac{1}{r} p_i \right) \end{aligned}$$

$$[B, E] = -\frac{\hbar^2 e^2}{m} \frac{r_i}{r^3}$$

$$[C, D] = -\frac{e^2}{2m} \left(\frac{1}{r} r_i p^2 - p^2 r_i \frac{1}{r} \right)$$

$$\begin{aligned}
 [P^2, \pi_i] &= p_j p_j \pi_i - \pi_i p_j p_j \\
 &= p_j \pi_i p_j - i\hbar p_i - \pi_i p_j^2 \\
 &= \pi_i p_j^2 - i\hbar p_i - \pi_i p_j^2 \\
 &= -2i\hbar p_i
 \end{aligned}$$

$$[C, D] = -\frac{e^2}{2m} \left(\frac{1}{\pi} \pi_i p^2 - \pi_i p^2 \frac{1}{\pi} + 2i\hbar p_i \frac{1}{\pi} \right)$$

$$\begin{aligned}
 [P^2, \frac{1}{\pi}] &= (p_i p_i \frac{1}{\pi} - \frac{1}{\pi} p_i p_i) \\
 &= (p_i \frac{1}{\pi} p_i + i\hbar p_i \pi_i \frac{1}{\pi^3} - \frac{1}{\pi} p_i p_i) \\
 &= i\hbar \frac{\pi_i}{\pi^3} p_i + i\hbar \pi_i p_i \frac{1}{\pi^3} + i\hbar (-i\hbar \times 3) \frac{1}{\pi^3} \\
 &= 2i\hbar \frac{\pi_i}{\pi^3} p_i + i\hbar \pi_i (3i\hbar \frac{\pi_i}{\pi^3}) + 3\hbar^2 \frac{1}{\pi^3} \\
 &= 2i\hbar \frac{1}{\pi^3} \vec{\pi} \cdot \vec{p}
 \end{aligned}$$

$$\begin{aligned}
 [C, D] &= -\frac{e^2}{2m} \left(\frac{1}{\pi} \pi_i p^2 - \pi_i \frac{1}{\pi} p^2 - \pi_i (2i\hbar \frac{1}{\pi^3} \vec{\pi} \cdot \vec{p}) \right. \\
 &\quad \left. + 2i\hbar \frac{1}{\pi} p_i + 2i\hbar (i\hbar \frac{\pi_i}{\pi^3}) \right)
 \end{aligned}$$

$$[C, D] = + \frac{e^2 \hbar}{2m} \left(i \frac{\pi_i}{\pi^3} \vec{\pi} \cdot \vec{p} - i \frac{1}{\pi} p_i + \hbar \frac{\pi_i}{\pi^3} \right)$$

$$\underline{[C, E] = 0}$$

$$[R_i, H] = \frac{e^2 \hbar}{m} \left(i \frac{4}{n} p_i - \frac{i \pi_i}{n^3} (\vec{n} \cdot \vec{p}) - \hbar \frac{\pi_i}{n^3} + i \frac{\pi_i}{n^3} \vec{n} \cdot \vec{p} - i \frac{1}{n} p_i + \hbar \frac{\pi_i}{n^3} \right)$$

= 0 !! OUF !!

⑤ R_i^2

$$[L_i, R_j] = \underbrace{\left[\hbar L_i, \frac{1}{2m} \sum_{k,l} \epsilon_{jkl} (p_k p_l - p_l p_k) \right]}_A + \underbrace{\left[L_i, \frac{\hbar^2}{n} \right]}_{i \hbar \epsilon_{ijk} \frac{\pi_k}{n}}$$

$$A = \frac{1}{2m} \epsilon_{jkl} (L_i p_k p_l - L_i p_l p_k - p_k p_l L_i + p_l p_k L_i)$$

$$= \frac{1}{2m} \epsilon_{jkl} (p_k p_l L_i + i \hbar \epsilon_{ikm} p_m p_l - p_l p_k L_i - i \hbar \epsilon_{ikm} p_m p_l - p_k p_l L_i + p_l p_k L_i)$$

$$= \frac{1}{2m} \epsilon_{jkl} (i \hbar \epsilon_{ikm} p_k p_m + i \hbar \epsilon_{ikm} p_m p_l - i \hbar \epsilon_{ikm} p_l p_m - i \hbar \epsilon_{ikm} p_m p_l)$$

$$= \frac{i \hbar}{2m} \underbrace{(\epsilon_{jkl} \epsilon_{ikm} + \epsilon_{ikm} \epsilon_{jkl})}_{B} p_k p_m - (\epsilon_{jkl} \epsilon_{ikm} + \epsilon_{ikm} \epsilon_{jlm}) p_l p_m$$

$$B = \cancel{S_{im} S_{il}} - \cancel{S_{ji} S_{jm}} + \cancel{S_{ij} S_{im}} - S_{je} S_{in}$$

$$= \epsilon_{ijk} \epsilon_{emk}$$

$$[R_i, R_j] =$$

$$A = \frac{i\hbar}{2m} \epsilon_{ijk} \epsilon_{klm} (p_l p_m - p_m p_l) \\ = i\hbar \epsilon_{ijk} \left(\frac{1}{2m} \epsilon_{klm} (p_l p_m - p_m p_l) \right)$$

$$\hookrightarrow [R_i, R_j] = i\hbar \epsilon_{ijk} R_k.$$

$$\text{[7]} \quad [p_i, p_j] = \frac{m}{-2E} [R_i, R_j] = i\hbar \epsilon_{ijk} p_k.$$

$$[L_i, p_j] = \sqrt{\frac{m}{-2E}} [L_i, R_j] = i\hbar \epsilon_{ijk} \sqrt{\frac{m}{-2E}} R_k \\ = i\hbar \epsilon_{ijk} p_k.$$

$$[L_i, L_j] = i\hbar \epsilon_{ijk} L_k$$

$$\text{[8]} \quad M \in \text{SO}(3) : \det M = +1$$

$${}^t M M = \mathbb{1}$$

$$M = e^{i \sum J_i \sigma_i} = \mathbb{1} + i \sum \sigma_i J_i$$

$${}^t M M = (\mathbb{1} + i \sum \sigma_i J_i) (\mathbb{1} + i \sum \sigma_j J_j)$$

$$= \mathbb{1} + i \sum \sigma_i ({}^t J_i + J_i) = \mathbb{1}$$

$\hookrightarrow J_i$ antisymétrique.

↳ $\frac{n(n-1)}{2}$ geraden

Sol 1) $\Rightarrow 1$

Sol 2) $\Rightarrow 3$

Sol 4) $\Rightarrow 6$

$$\begin{aligned}
 [J_{mn}, J_{pq}] &= (J_{mn})_{ik} (J_{pq})_{kj} - (J_{pq})_{ki} (J_{mn})_{kj} \\
 &= - (\delta_{mi} \delta_{nk} - \delta_{mk} \delta_{ni}) (\delta_{pk} \delta_{qj} - \delta_{pj} \delta_{qk}) + (\delta_{pi} \delta_{qk} - \delta_{pk} \delta_{qi}) (\delta_{mj} \delta_{nk} - \delta_{mk} \delta_{nj}) \\
 &= - (\delta_{mi} \delta_{np} \delta_{qj} - \delta_{mi} \delta_{nq} \delta_{pj} - \delta_{mp} \delta_{ni} \delta_{qj} + \delta_{mq} \delta_{ni} \delta_{pj}) \\
 &\quad + (\delta_{pi} \delta_{qn} \delta_{nj} - \delta_{pi} \delta_{qn} \delta_{mj} - \delta_{pn} \delta_{qi} \delta_{nj} + \delta_{pn} \delta_{qi} \delta_{mj}) \\
 &= \delta_{np} (\delta_{qi} \delta_{mj} - \delta_{mi} \delta_{qj}) + \delta_{mq} (\delta_{pi} \delta_{nj} - \delta_{ni} \delta_{pj}) \\
 &\quad - \delta_{np} (-\delta_{ni} \delta_{qj} + \delta_{qj} \delta_{ni}) - \delta_{mq} (\delta_{pi} \delta_{mj} - \delta_{mi} \delta_{pj}) \\
 &= i \{ \delta_{np} J_{qn} + \delta_{nq} J_{pn} - \delta_{mp} J_{qn} - \delta_{mq} J_{pn} \} \\
 &= i (\delta_{np} J_{nq} + \delta_{nq} J_{np} - \delta_{mp} J_{np} - \delta_{mq} J_{mq})
 \end{aligned}$$

Q) $[a_i, a_j] = [J_{0i}, J_{0j}] = i (J_{ij}) = i \epsilon_{ijk} a_k$

$$\begin{aligned}
 [b_i, a_j] &= \frac{1}{2} \epsilon_{ike} [J_{ke}, J_{0j}] \\
 &= \frac{i}{2} \epsilon_{ike} (-\delta_{kj} J_{0e} + \delta_{ej} J_{0k}) \\
 &= \frac{i}{2} (+\epsilon_{ije} a_e - \epsilon_{ikj} a_k) \\
 &= i \epsilon_{ijk} a_k
 \end{aligned}$$

$$\begin{aligned}
 [L_i, L_j] &= \frac{1}{4} \epsilon_{ilm} \epsilon_{jpn} [J_{em}, J_{pn}] \\
 &= \frac{i}{4} \epsilon_{ilm} \epsilon_{jpn} \{ \delta_{pp} J_{mq} + \delta_{mq} J_{ep} - \delta_{eq} J_{mp} - \delta_{mp} J_{eq} \} \\
 &= \frac{i}{4} \{ -J_{ji} - J_{ji} - J_{ji} - J_{ji} \} = i J_{ij} \\
 &= i \epsilon_{ijk} L_k.
 \end{aligned}$$

\vec{a} joue le rôle que \vec{k}/\hbar } L'algèbre d'algèbre présente une
 \vec{b} " " " " \vec{e}/\hbar } symétrie $SO(4)$! beaucoup plus
 grand que $SO(3)$...

$$\begin{aligned}
 10) [M_i, M_j] &= \frac{1}{4} ([L_i, L_j] + [L_i, K_j] + [K_i, L_j] + [K_i, K_j]) \\
 &= \frac{i\hbar}{4} \epsilon_{ijk} (L_k + K_k + K_k + L_k) \\
 &= i\hbar \epsilon_{ijk} M_k.
 \end{aligned}$$

$$\begin{aligned}
 [M_i, N_j] &= \frac{1}{4} ([L_i, L_j] - [L_i, K_j] + [K_i, L_j] - [K_i, K_j]) \\
 &= \frac{i\hbar}{4} \epsilon_{ijk} (L_k - K_k + K_k - L_k) = 0
 \end{aligned}$$

$$\begin{aligned}
 [N_i, N_j] &= \frac{i\hbar}{4} \epsilon_{ijk} (L_k - K_k - K_k + L_k) \\
 &= i\hbar \epsilon_{ijk} N_k.
 \end{aligned}$$

\hookrightarrow Les M_i et les N_j commutent entre eux, et forment chacun l'algèbre de $SO(3)$.

$$\begin{aligned}
 \hookrightarrow SO(4) &\sim SU(2) \times SU(2) \\
 &\sim SO(3) \times SO(3).
 \end{aligned}$$

$$M^2 = \frac{1}{4} (p^2 + h^2 + \vec{p} \cdot \vec{h} + \vec{h} \cdot \vec{p})$$

$$N^2 = \frac{1}{4} (p^2 + h^2 - \vec{p} \cdot \vec{h} - \vec{h} \cdot \vec{p})$$

$$= \frac{1}{4} (p^2 + h^2) = M^2$$

na rechten proprie de M^2 : $M^2 |nm\rangle = N^2 |nm\rangle$

\rightarrow rechten proprie de N^2 .

12)

$$R^2 = -\frac{m}{2E} R^2$$

~~$R^2 + E^2$~~

$$R^2 + E^2 = 4M^2$$

$$-\frac{m}{2E} \left(e^4 + \frac{2E(p^2 + h^2)}{m} \right) + E^2 = 4M^2$$

$$-\frac{m e^4}{2E} - h^2 = 4M^2$$

$$-\frac{m e^4}{2E} = h^2 (4h(h+1) + 1)$$

$$= h^2 (2h+1)^2$$

$$E = -\frac{m e^4}{2h^2 (2h+1)^2} = -\frac{R_H}{(2h+1)^2} \quad \text{or} \quad R_H = \frac{m e^4}{2h^2}$$

1) \times Loi interne.

\times identité: $\{123\} \times g = g \times \{123\} = g$.

\times associativité OK.

\times inverse: OK.

\times

- il y a 6 éléments

$\{123\}$
 $\{132\}$

$\{213\}$
 $\{231\}$

$\{312\}$
 $\{321\}$.

111

~~111~~

41

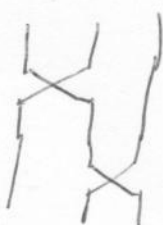
~~41~~

~~111~~

~~111~~

~~111~~

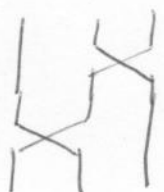
- groupe non commutatif:



=



{



=



2) $\{123\}: A_1 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$

$\{213\}: A_2 = \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$

$\{312\}: A_3 = \begin{pmatrix} 0 & 0 & 1 \\ 1 & 0 & 0 \\ 0 & 1 & 0 \end{pmatrix}$

$\{132\}: A_4 = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}$

$\{231\}: A_5 = \begin{pmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}$

$\{321\}: A_6 = \begin{pmatrix} 0 & 0 & 1 \\ 0 & 1 & 0 \\ 1 & 0 & 0 \end{pmatrix}$

$A_i \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$

\hookrightarrow Le vecteur $\begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$ ne se transforme pas: représentation triviale fermée

3) on introduit $M_1 = \frac{1}{\sqrt{3}} \begin{pmatrix} 1 \\ 1 \\ 1 \end{pmatrix}$

$M_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -1 \\ 0 \end{pmatrix}$

$M_3 = \sqrt{\frac{2}{3}} \begin{pmatrix} 1/2 \\ +1/2 \\ -1 \end{pmatrix}$

o vérifie facilement que $\vec{M}_i \cdot \vec{M}_j = \delta_{ij}$

o construit $S = (\vec{M}_1, \vec{M}_2, \vec{M}_3)$

$$A'_i = {}^t S A_i S.$$

$$(A'_i)_{ii} = 1$$

$(A'_i)_{ia}$ et $(A'_i)_{a,i}$ $a=2,3$ sont obtenus comme un produit scalaire

entre \vec{M}_i avec les coordonnées permises et \vec{M}_2, \vec{M}_3
 $= \vec{M}_1$

↳ $(A'_i)_{ia} = (A'_i)_{a,i} = 0 \Rightarrow$ matrice bloc diagonale

↳ Les 6 matrices correspondent aux rotations le long de $(1,1,1)$

d'angle $\pm \frac{2\pi}{3}$

et aux images miroir / ex plus $(x0z), (y0z)$ et $(z0x)$.

ces opérations laissent (111) inchangé.

Les ~~axes~~ Les parties orthogonale à cette direction se mesurent
entre elles.

$$1) T'_{ii} = R_{ik} R_{ie} T_{ke}$$

$$R \in SO(3) \Rightarrow {}^t R R = \mathbb{1}$$

$$T'_{ii} = S_{ke} T_{ke} = T_{kk}$$

↳ la trace est invariante.

↳ la transformée de la trace ne dépend que de la trace \Rightarrow fait intervenir
représentation (triviale) du groupe de rotation.

$$\begin{aligned} 2) A'_{ij} &= \frac{1}{2} (R_{ik} R_{je} T_{ke} - R_{jk} R_{ie} T_{ke}) \\ &= \frac{1}{2} R_{ik} R_{je} (T_{ke} - T_{ek}) \\ &= R_{ik} R_{je} A_{ke} \end{aligned}$$

La transformée de A ne fait intervenir que A \Rightarrow représentation invariante

$$V'_i = \frac{1}{2} \epsilon_{ijk} R_{je} R_{km} A_{em}$$

$$\begin{aligned} \epsilon_{jle} R_{je} R_{km} &= \epsilon_{ijk} (S_{je} + i\sigma_2 J_{je}^d) (S_{km} + i\sigma_2 J_{km}^d) \\ &= \epsilon_{ilm} + i\sigma_2 (\epsilon_{iek} (-i) \epsilon_{p \times km} + \epsilon_{ijm} (-i) \epsilon_{ajl}) \\ &\quad + o(\sigma^2) \\ &= \epsilon_{ilm} + i\sigma_2 (-i) (\sin \sigma_{ek} - S_{ik} S_{en} + S_{ia} S_{me} - S_{ie} S_{na}) \\ &= \epsilon_{ilm} + i\sigma_2 (-i) (\epsilon_{iak} \epsilon_{m \times lk}) \end{aligned}$$

$$V'_i = V_i + i \sigma_a (-i) \epsilon_{ikl} V_k$$

↳ se transforme comme un vecteur \Rightarrow

A: représentation vectorielle du groupe des rotations

$$\begin{aligned} 3] S'_{ij} &= \frac{1}{2} (T'_{ij} + T'_{ji}) = \frac{1}{2} (R_{ik} R_{je} T_{ke} + R_{jk} R_{ie} T_{ke}) \\ &= R_{ik} R_{je} S_{ke} \end{aligned}$$

4] générateurs des rotations autour de x :

$$S'_{ij} = S_{ij} + i \theta (-i) (S_{ik} \epsilon_{ljk} + S_{jk} \epsilon_{lik}) S_{ke}$$

$$S'_{11} = (0)$$

$$S'_{12} = (S_{13})$$

$$S'_{13} = (-S_{12})$$

$$S'_{22} = (2S_{23})$$

$$S'_{23} = (-S_{22} + S_{33})$$

$$S'_{33} = (-2S_{23})$$

$$J_x = -i \begin{pmatrix} 0 & 0 & 0 & 0 & 2 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 4 \\ 0 & 0 & 0 & -1 & 0 \end{pmatrix}$$

relaciones entre de y:

$$\begin{aligned} S'_{22} &\rightarrow 0 \\ S'_{23} &= S_{12} \\ S'_{12} &= -S_{23} \\ S'_{33} &= 2S_{13} \\ S'_{13} &= S_{11} - S_{33} \\ S'_{11} &= -2S_{13} \end{aligned}$$

$$J_y = -i \begin{pmatrix} 0 & 0 & -4 & 0 & 0 \\ 0 & 0 & 0 & 0 & -1 \\ 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & -2 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \end{pmatrix}$$

relaciones entre de z:

$$\begin{aligned} S'_{33} &= 0 \\ S'_{13} &= S_{23} \\ S'_{23} &= -S_{13} \\ S'_{11} &= 2S_{12} \\ S'_{12} &= S_{22} - S_{11} \\ S'_{22} &= -2S_{12} \end{aligned}$$

$$J_z = -i \begin{pmatrix} 0 & 2 & 0 & 0 & 0 \\ -1 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & 1 \\ 0 & -2 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 \end{pmatrix}$$

On retrouve les relations de commutation...

$[R_i, R_j]$? on calcule $[R_x, R_y]$

ANNEXE:
Calcul des relations
supplémentaires...

$$R_x = \frac{1}{m} (P_y l_z - P_z l_y - i\hbar P_x) - e^2 \frac{x}{r}$$

$$R_y = \frac{1}{m} (P_z l_x - P_x l_z - i\hbar P_y) - e^2 \frac{y}{r}$$

$$[R_x, R_y] = \frac{1}{m^2} [P_y l_z - P_z l_y - i\hbar P_x, P_z l_x - P_x l_z - i\hbar P_y] \\ - \frac{e^2}{m} \left\{ [P_y l_z - P_z l_y - i\hbar P_x, \frac{y}{r}] - [P_z l_x - P_x l_z - i\hbar P_y, \frac{x}{r}] \right\}$$

$$[P_y l_z, P_z l_x] = P_y l_z P_z l_x - P_z l_x P_y l_z \\ = P_z P_y l_x l_z + i\hbar P_z P_y l_y - P_z l_x P_y l_z \\ = -i\hbar P_z^2 l_z + i\hbar P_z P_y l_y$$

$$[P_y l_z, P_x l_z] = P_y l_z P_x l_z - P_x l_z P_y l_z \\ = P_y P_x l_z l_z + i\hbar P_y^2 l_z - P_x l_z P_y l_z \\ = i\hbar P_x^2 l_z + i\hbar P_y^2 l_z$$

~~$[P_z l_y, P_x l_z] =$~~

$$[P_z l_y, P_z l_x] = P_z l_y P_z l_x - P_z l_x P_z l_y \\ = i\hbar P_z l_y P_y + P_z l_y l_x P_z - P_z l_x P_z l_y \\ = i\hbar P_z l_y P_y + P_z l_x l_y P_z - i\hbar P_z l_z P_z - P_z l_x P_z l_y \\ = i\hbar (P_z P_y l_y + P_x P_z l_x - P_z^2 l_z)$$

$$[P_z l_y, P_x l_z] = P_z l_y P_x l_z - P_x l_z P_z l_y \\ = P_z P_x l_y l_z - i\hbar P_z^2 l_z - P_x l_z P_z l_y \\ = P_z P_x l_y l_z + i\hbar P_z P_x l_x - i\hbar P_z^2 l_z - P_x l_z P_z l_y \\ = i\hbar (P_z P_x l_x - P_z^2 l_z)$$

Le premier commutateur

$$i\hbar \left\{ \begin{aligned} & -P_z^2 L_z + P_y P_z L_y - P_x^2 L_z - P_y^2 L_z - P_y P_z L_y - P_x P_z L_x + P_z^2 L_z \\ & + P_x P_z L_x - P_z^2 L_z - P_x P_z L_x + P_z L_x P_x + P_x^2 L_z - P_x L_z P_x \\ & - P_y L_z P_y + P_y^2 L_z + P_z P_y L_y - P_y P_z L_y \end{aligned} \right\}$$

$$= i\hbar \left\{ -P_z^2 L_z - P_x^2 L_z - i\hbar P_x P_y + i\hbar P_y P_x - P_y^2 L_z \right\}$$

$$= -i\hbar P^2 L_z$$

$$[P_x, x^2] = P_x x x - x x P_x$$

$$= -i\hbar x + x P_x x - x x P_x$$

$$= -2i\hbar x$$

$$[P_x, (x^2)^{n-1}] = -2i\hbar x (x^2)^{n-1}$$

$$[P_x, f(x)] = -2i\hbar x f'(x)$$

$$[P_x, \frac{1}{n}] = i\hbar \frac{x}{n^2}$$

$$[P_y L_z, \frac{y}{n}] = P_y L_z y \frac{1}{n} - y \frac{1}{n} P_y L_z$$

$$= P_y y \frac{1}{n} L_z - i\hbar P_y x \frac{1}{n} - y \frac{1}{n} P_y L_z$$

$$= y P_y \frac{1}{n} L_z - i\hbar P_y \frac{x}{n} - y \frac{1}{n} P_y L_z - i\hbar \frac{1}{n} L_z$$

$$= \frac{iy^2 \hbar}{n^2} L_z - i\hbar \frac{1}{n} L_z - i\hbar P_y \frac{x}{n}$$

$$[P_z L_y, \frac{y}{n}] = P_z L_y y \frac{1}{n} - y \frac{1}{n} P_z L_y$$

$$= y \frac{i\hbar z}{n^2} L_y$$

$$[P_z L_x, \frac{x}{n}] = P_z L_x x \frac{1}{n} - x \frac{1}{n} P_z L_x$$

$$= i\hbar \frac{x z}{n^2} L_x$$

$$[P_x L_z, \frac{x}{n}] = P_x L_z x \frac{1}{n} - x \frac{1}{n} P_x L_z$$

$$= P_x x L_z \frac{1}{n} + i\hbar P_x y \frac{1}{n} - x \frac{1}{n} P_x L_z$$

$$= x P_x L_z \frac{1}{n} + i\hbar L_z \frac{1}{n} + i\hbar P_x y \frac{1}{n} - x \frac{1}{n} P_x L_z$$

$$= i\hbar \frac{x^2}{n^2} L_z - i\hbar \frac{1}{n} L_z + i\hbar P_x y \frac{1}{n}$$

$2e^e + 3e^e$ comutatoras

(c)

$$\frac{i\hbar}{n} \left\{ \frac{y^2}{n^2} p_z - p_z - n p_x \frac{x}{n} - \frac{yz}{n^2} p_y - \frac{xz}{n^2} p_x + \frac{x^2}{n^2} p_z - p_z + n p_x \frac{y}{n} \right\}$$

$$+ i\hbar \left\{ - p_x \frac{y}{n} + \frac{y}{n} p_x + p_x \frac{x}{n} - \frac{x}{n} p_x \right\}$$

$$yzp_y + xzp_x = yz^2 p_x - yz^2 p_x + xz^2 p_y - xz^2 p_y$$

$$= -z^2 p_z$$

$$\hookrightarrow \frac{i\hbar}{n} \{ p_z - p_z - p_z - p_z \} = -\frac{2i\hbar}{n} p_z$$

$$\hookrightarrow [R_x, R_y] = -\frac{i\hbar p^2}{m^2} p_z + \frac{e^2}{m} \frac{2i\hbar}{n} p_z$$

$$= -\frac{2i\hbar}{m} \left(\frac{p^2}{2m} - \frac{e^2}{n} \right) p_z$$

$$= -\frac{2i\hbar}{m} H p_z$$

$$\hookrightarrow [R_i, R_j] = -\frac{2i\hbar}{m} H \epsilon_{ijk} p_k \quad \text{OUF!}$$

$$R_x = \frac{1}{m} (P_y l_z - P_z l_y) - \frac{i\hbar}{m} P_x - e^2 \frac{x}{r}$$

$$R_x^2 = \frac{1}{m^2} (P_y l_z - P_z l_y)(P_y l_z - P_z l_y) - \frac{\hbar^2}{m^2} P_x^2 + e^4 \frac{x^2}{r^2} \\ - \frac{i\hbar}{m^2} \left\{ P_x, P_y l_z - P_z l_y \right\} - \frac{e^2}{m} \left\{ P_y l_z - P_z l_y, \frac{x}{r} \right\} \\ + \frac{i\hbar e^2}{m} \left\{ P_x, \frac{x}{r} \right\}$$

$$A = P_y l_z P_y l_z + P_z l_y P_z l_y - P_y l_z P_z l_y - P_z l_y P_y l_z \\ = P_y^2 l_z^2 + P_z^2 l_y^2 - i\hbar P_y P_x l_z + i\hbar P_z P_x l_y - \underbrace{\left\{ P_y (x l_z - y l_x), P_z (z l_x - y l_y) \right\}}_E$$

$$E = \underbrace{P_y x P_y P_z z P_x}_a + \underbrace{P_z z P_x P_y x P_y}_a + \underbrace{P_y y P_x P_z x P_z}_b + \underbrace{P_z x P_z P_y y P_x}_b \\ - \underbrace{P_y x P_y P_z x P_z}_a - \underbrace{P_z x P_z P_y x P_y}_b - \underbrace{P_y y P_x P_z z P_x}_c - \underbrace{P_z z P_z P_y y P_x}_c \\ - \underbrace{P_y^2 z P_x z P_x}_a + \underbrace{P_z^2 x P_x x P_x}_c - \underbrace{P_z^2 y P_x y P_x}_b + \underbrace{P_z^2 y P_x y P_x}_c$$

$$= -P_y^2 (x P_x z P_x + z P_x z P_x - x P_z z P_x - z P_x x P_z + i\hbar P_x x) \\ - P_z^2 (x P_y x P_y + y P_x y P_x - x P_y y P_x - y P_x x P_y + i\hbar P_x x) \\ + P_x^2 (z P_y z P_y + y P_z y P_z - y P_z z P_y + i\hbar P_z z - z P_y y P_z + i\hbar P_y y) \\ = -P_y^2 l_y^2 - P_z^2 l_z^2 + P_x^2 l_x^2 + i\hbar (-P_y^2 P_x x - P_z^2 P_x x + P_x^2 P_z z + P_x^2 P_y y)$$

$$A = P_y^2 l_z^2 + P_z^2 l_y^2 + P_y^2 l_y^2 + P_z^2 l_z^2 - P_x^2 l_x^2 + i\hbar \left\{ P_x P_y z P_x - P_x P_z x P_x \right. \\ \left. - P_x P_y x P_y + P_x P_z y P_x + P_y^2 P_x x + P_z^2 P_x x - P_x^2 P_z z - P_x^2 P_y y \right\}$$

$$A = (P_y^2 + P_z^2)(P_y^2 + P_z^2) - P_x^2 P_x^2$$

$$\begin{aligned} B &= P_x P_y x P_y + P_y x P_y P_x - P_x P_y y P_x - P_y y P_x P_x \\ &\quad - P_x P_z z P_x - P_z z P_x P_x + P_x P_z x P_z + P_z x P_z P_x \\ &= 2 P_y^2 P_x x + i \hbar P_y^2 - 2 P_x^2 P_y y - 2 P_x^2 P_z z + 2 P_z^2 P_x x + i \hbar P_z^2 \\ &= 2 \left(\frac{P_y^2 + P_z^2}{\hbar} P_x x - P_x^2 (P_y y + P_z z) \right) + i \hbar (P_y^2 + P_z^2) \\ &\quad \sum R_i^2 \Rightarrow 0. \end{aligned}$$

$$\begin{aligned} [P_x, x^2] &= P_x x^2 - x^2 P_x = -i \hbar x + x P_x x - x^2 P_x \\ &= -2i \hbar x \end{aligned}$$

$$[P_x, f(x)] = -i \hbar x f'(x)$$

$$[P_x, \frac{1}{x}] = i \hbar \frac{x}{x^3}$$

$$\begin{aligned} C &= P_y p_z \frac{x}{\hbar} + \frac{x}{\hbar} P_y p_z - P_z p_y \frac{x}{\hbar} - \frac{x}{\hbar} P_z p_y \\ &= x P_y \frac{1}{\hbar} p_z + i \hbar P_y \frac{y}{\hbar} + \frac{x}{\hbar} P_y p_z - x P_z \frac{1}{\hbar} p_y + i \hbar P_z \frac{z}{\hbar} - \frac{x}{\hbar} P_z p_y \\ &= \frac{2}{\hbar} \left(x P_y p_z - x P_z p_y \right) + i \hbar \left(\frac{x y}{\hbar^2} p_z + P_y \frac{y}{\hbar} - \frac{x z}{\hbar^2} p_y + P_z \frac{z}{\hbar} \right) \end{aligned}$$

$$\frac{-e^2}{m} \left(+i \hbar \frac{e^2}{\hbar} D \right)$$

$$\begin{aligned} &= \frac{e^2}{m} \left\{ -\frac{2}{\hbar} (x P_y p_z - x P_z p_y) + i \hbar \left(-\frac{x y}{\hbar^2} p_z + \frac{x z}{\hbar^2} p_y - P_y \frac{y}{\hbar} + P_z \frac{z}{\hbar} \right. \right. \\ &\quad \left. \left. + P_x \frac{x}{\hbar} + \frac{x}{\hbar} P_x \right) \right\} \end{aligned}$$

$$\begin{aligned} \frac{x}{\hbar} P_x &= \frac{1}{\hbar} P_x x + i \hbar \frac{1}{\hbar} \\ &= P_x \frac{x}{\hbar} - i \hbar \frac{x^2}{\hbar^2} + i \hbar \frac{1}{\hbar} \end{aligned}$$

$$-\frac{e^2}{m} C + \frac{i\hbar}{m} e^2 D = \frac{e^2}{m} \left\{ -\frac{2}{\pi} (x P_y l_z - x P_z l_y) + i\hbar \left(-\frac{x y}{n^3} P_z + \frac{x z}{n^3} l_y \right) \right. \\ \left. - \underbrace{\left(P_y \frac{y}{n} - P_z \frac{z}{n} + P_x \frac{x}{n} \right)}_{\sum_i R_i^2 = 0} - \hbar^2 \left(\frac{1}{n} - \frac{2l^2}{n^3} \right) \right\}$$

$$\sum_i R_i^2 = \frac{1}{m^2} P^2 e^2 - \frac{\hbar^2}{m^2} P^2 + e^4 + \frac{\hbar^2}{m^2} 2P^2 + \frac{e^2}{m} \left\{ -\frac{2}{n} e^2 - \hbar^2 \frac{2}{n} \right\} \\ = \frac{2}{m} \left\{ \frac{1}{2m} P^2 (e^2 + \hbar^2) - \frac{e^2}{n} (e^2 + \hbar^2) \right\} - e^4$$

$$R^2 = \frac{2\hbar}{m} (P^2 + \hbar^2) + e^4 \quad \text{OUF!}$$